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Китик Иван, Мирончук Галина, Парасюк Олег, Данильчук Сергей, Божко Владимир, Замуруева Оксана. Оптические свойства кристаллов системы $Tl_{1-x}In_xSn_xS_2$ ($x=0; 0,1; 0,20; 0,3; 0,4; 0,5$). В данной работе изучалась кристаллическая система $Tl_{1-x}In_xSn_xS_2$ ($x = 0,1; 0,20; 0,3; 0,4; 0,5$), а именно влияние частичного катионного замещения ионов In ионами Sn на оптические характеристики исследуемых соединений. Рассмотрены закономерности изменения оптических спектров поглощения в кристаллах твёрдых растворов от температуры и состава. Оценена ширина запрещенной зоны соединений для различных температур, прямых и непрямых переходов. Определены по экспериментальной зависимости $\ln\alpha(h\nu)$ при различных температурах $E_g(0)$ и α_0 . Рассчитана концентрация заряженных дефектов, ответственных за экспериментально установленное значение Δ_0 .

Ключевые слова: коэффициент поглощения, ширина запрещенной зоны, слоистые структуры.

Kituk Ivan, Myronchuk Galina, Parasyuk Oleg, Danylchuk Sergej, Bojko Volodymyr, Zamuruyeva Oksana. Optical properties of crystals of $Tl_{1-x}In_xSn_xS_2$ ($x=0; 0,1; 0,20; 0,3; 0,4; 0,5$). The influence of partial cationic substitution of In ions by Sn ions on the optical characteristics of the solid solutions $Tl_{1-x}In_xSn_xS_2$ ($x = 0,1; 0,20; 0,3; 0,4; 0,5$) has been studied. A comparative analysis of temperature behavior of optical absorption edge over the temperature range from 77 to 300 K is performed for alloys with ($x = 0,1; 0,20; 0,3; 0,4; 0,5$).

The band gap of compounds at different temperatures, for direct and indirect transitions have been estimated. From the experimental dependences $\ln\alpha(h\nu)$ at different temperatures the constants $E_g(0)$ and α_0 in the Urbach rule have been determined. The concentration of charged defects responsible for the experimentally determined value of characteristic energy has been calculated.

Key words: absorption, band gap, layered structure.

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Electronic structure of $AgCd_2GaS_4$

In the present work we report on measurements of *X-ray* photoelectron (*XP*) valence-band spectrum including the comparatively wide energy region corresponding to location of upper core-levels (up to ~ 80 eV with respect to the bottom of the valence band) of the atoms constituting a quaternary $AgCd_2GaS_4$ single crystal grown by the Bridgman

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method. Electronic structure of $AgCd_2GaS_4$ are also theoretically studied by using the first-principles full potential linearized augmented plane wave (FP-LAPW) method. The FP-LAPW method has been employed in the present work to calculate binding energies of upper core-levels for the $AgCd_2GaS_4$ compound. With respect to values of binding energies of the Cd 4p, Ag 4p and Ga 3d core-levels, the theoretical and experimental data for the $AgCd_2GaS_4$ compound were found to be in agreement to each other. Curves of dominant partial densities of states of $AgCd_2GaS_4$ have been compared on a common energy scale with the XP valence-band spectrum of the compound under consideration.

Key words: $AgCd_2GaS_4$; Electronic structure; First-principles band-structure calculations; X-ray photoelectron spectroscopy; densities of states.

Introduction. A quaternary $AgCd_2GaS_4$ compound crystallizes in the orthorhombic structure (space group $Pmn2_1$; lattice parameters: $a = 8.1395(9) \text{ \AA}$, $b = 6.9394(8) \text{ \AA}$, and $c = 6.6014(7) \text{ \AA}$ [11]). It is formed in accordance with a peritectic reaction $L + \beta(CdS) \rightleftharpoons AgCd_2GaS_4$ and exists in the 67–78 mol % CdS homogeneity region [12]. Tetrahedral coordination of the constituent atoms is characteristic of the unit cell of $AgCd_2GaS_4$ [15]. In the structure of $AgCd_2GaS_4$, there are three non-equivalent types of sulfur atoms (Fig. 1), which occupy $4b$ and two different $2a$ sites. The non-central space group of the crystal structure of the $AgCd_2GaS_4$ compound makes it a very prospective material for optoelectronics and non-linear optics. The $AgCd_2GaS_4$ compound reveals intense luminescence, high photosensitivity, and wide transparency range [13,14]. The compound is n-type semiconductor with the width of energy gap, E_g , of 2.15 eV at about 300 K.

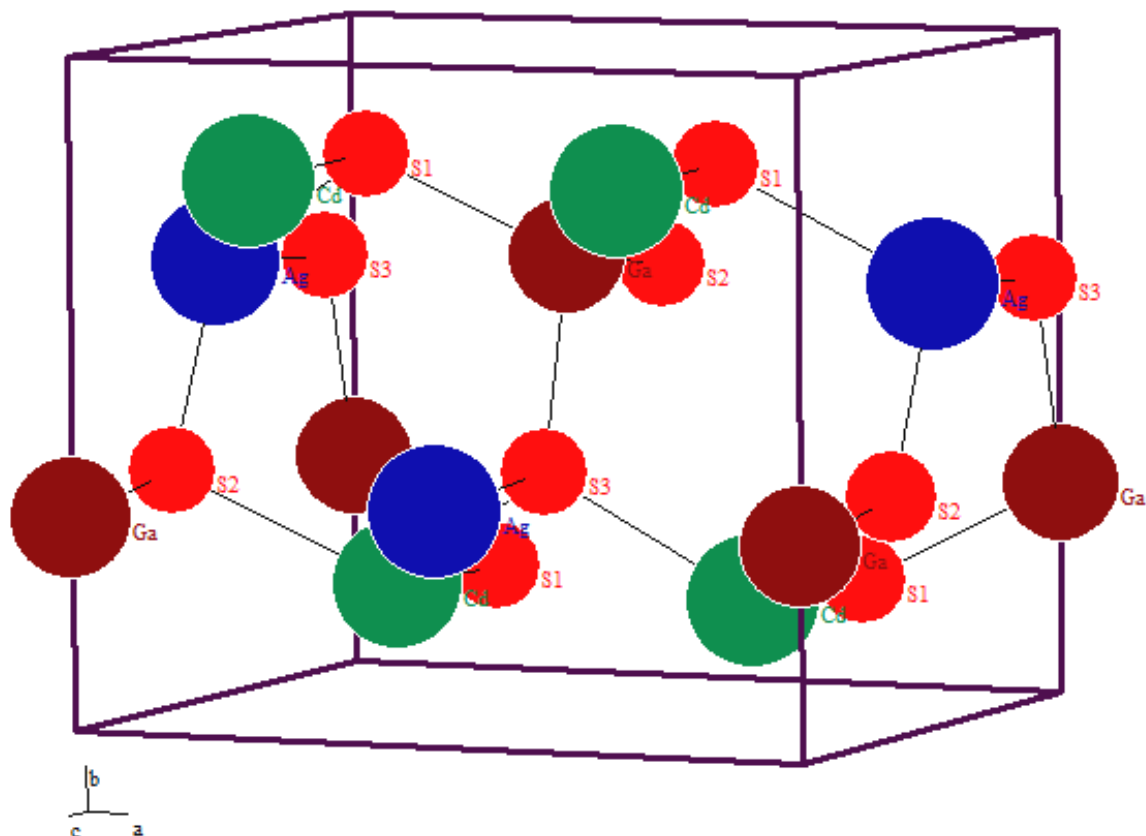


Fig. 1. Crystal structure of $AgCd_2GaS_4$ (unit cell is outlined)

The valence-band occupations and optical susceptibilities of $AgCd_2GaS_4$ have been calculated comparatively recently by Reshak et al. [4; 8]. It has been established that the valence band maximum and conduction band minimum are located at the Γ point of the Brillouin zone. This means that $AgCd_2GaS_4$ is a direct-gap semiconductor. In addition, Reshak et al. [4,8] have established a strong negative uniaxial anisotropy for optical properties of $AgCd_2GaS_4$. Recently in Ref. [7], X-ray photoelectron (XP) core-level and valence-band spectra for pristine and Ar^+ -ion irradiated (001) surface of $AgCd_2GaS_4$ single crystal have been measured and those measurements have revealed high chemical stability of (001) $AgCd_2GaS_4$ surface.

Further, electronic structure of $AgCd_2GaS_4$ has been calculated employing the full potential linearized augmented plane wave (FP-LAPW) method [7]. Additionally, in Ref. [7]. X-ray emission bands representing the energy distribution of the valence $Ag\ d$, $Cd\ d$, $Ga\ p$ and $S\ p$ states have been recorded and compared on a common energy scale with the XP valence-band spectrum of the $AgCd_2GaS_4$ compound. The FP-LAPW calculations of the electronic structure of the $AgCd_2GaS_4$ compound have revealed that its upper valence band is composed mainly from contributions of the $Ag\ d$ states, with also significant contributions of the p states associated with S and Ga atoms, whilst the bottom of the conduction band of $AgCd_2GaS_4$ is dominated by contributions of the unoccupied $Ga\ s$ states [7].

In the present paper we aim at a study of XP valence-band spectrum including the comparatively wide energy region corresponding to location of upper core-levels of the atoms constituting a quaternary $AgCd_2GaS_4$ compound (up to $\sim 80\ eV$ with respect to the bottom of the valence band). Additionally, we aim at first-principles calculations of the electronic structure of $AgCd_2GaS_4$. Therefore, ab initio FP-LAPW method has been employed in the present work to calculate binding energies of upper core-levels of the $AgCd_2GaS_4$ compound. Curves of dominant partial densities of states of $AgCd_2GaS_4$ have been also compared in the present work on a common energy scale with the XP valence-band spectrum of the compound under consideration.

Basic results and its discussions

1. Experimental. In the present study, we have used the $AgCd_2GaS_4$ single crystal grown by the Bridgman method. Its synthesis was described in detail in [7]. Briefly, for the growth, a 20 g batch of appropriately weighed amounts of high-purity elementary substances (purity of Ag, Cd, Ga and S better than 99,999 wt.%) was placed into an ampoule, following by its evacuating and soldering. We have used the preliminary synthesis of the batch by the direct bonding of the elementary sulfur by heating the ampoule in the flame of the oxygen-gas burner and further homogenization in the rotating furnace. The actual growth process was performed in a vertical two-zone furnace by the gradual growth of the crystal onto the formed seed. A high-quality $AgCd_2GaS_4$ single crystal was derived using the following growth conditions: growth zone/annealing zone temperature of the furnace $\sim 1100/650^\circ C$; temperature gradient at the solid-melt interface $\sim 2,0^\circ/mm$; growth rate 0.15 mm/h; annealing duration 100 h. The diffraction pattern of the compound was indexed in the space group $Pmn2_1$, and the refined unit cell periods being in agreement with the data originally reported for this compound in Refs [11,14].

The $AgCd_2GaS_4$ single crystal for the present XPS measurements was prepared in the shape of a thin plate with the following dimensions: length $\sim 7\ mm$, width $\sim 5\ mm$, and height $\sim 1\ mm$. The plate was cut from the boule to obtain (001) surface. The plate was polished in castor oil containing additions of beryllium abrasive No. 28 to gain a high-quality optical surface following the technique reported in Ref [5]. Fig. 2 presents a photo image of the $AgCd_2GaS_4$ plate used in the present experimental studies.

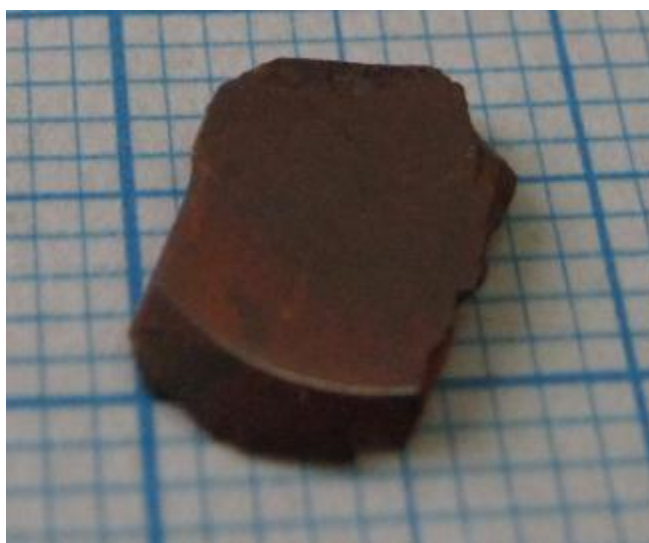


Fig. 2. Panoramic view of the polished $AgCd_2GaS_4$ plate

XPS spectra of the $AgCd_2GaS_4$ single crystal were recorded in an ion-pumped chamber having a base pressure less than 5×10^{-10} mbar of the UHV-Analysis-System assembled by SPECS Company (Germany). The spectra were excited employing a $Mg K\alpha$ source of X-ray radiation ($E=1253.6$ eV). The energy scale of the spectrometer was calibrated as reported in Refs. [5,7]. The charging effect was taken into account in reference to the C 1s line (284.6 eV) of adventitious carbon.

2. Method of calculation. Calculations of electronic properties of $AgCd_2GaS_4$ were fulfilled employing the first-principles self-consistent *FP-LAPW* method with the WIEN97 code [1]. We have used the lattice parameters and positions of the constituent atoms of $AgCd_2GaS_4$ determined for the crystal in Ref [13]. The *muffin-tin* sphere radii of the atoms constituting $AgCd_2GaS_4$ in the present calculations have been chosen as following: 1.376 Å for Ag and Cd, 1.164 Å for Ga, and 1.058 Å for S. The Rk_{max} parameter determining a number of basis functions equals 7.0 ($G_{MAX} = 14$). The order of the secular matrix was ~ 1800 – 1900 for different k -points of the Brillouin zone. Non-spherical contributions up to $l = 4$ have been taken into account in the potential decomposition. The generalized gradient approximation by Perdew et al. [10] was used for calculations of the exchange-correlation potential. The basis function consists of the same atomic orbitals of Ag, Cd, Ga, and S as reported in Ref [7]. Integration through the Brillouin zone was done using the tetrahedron method by Blöchl et al. [2]. The iteration process was checked taking into account changes of total energy and the calculations were interrupted when for three following iterations the change of total energy was less than 1.8×10^{-6} eV.

3. Results and discussion. Fig. 3 presents XP valence-band spectrum including the comparatively wide energy region corresponding to location of upper core-levels of the atoms constituting a quaternary $AgCd_2GaS_4$ compound (up to ~ 80 eV with respect to the bottom of the valence band). From this figure, it is obvious that the XP Cd 4p, Ag 4p, Ga 3d, and Cd 4d core-level spectra binding energies have their maxima at about 67.1, 57.2, 19.6, and 10.8 eV, respectively. To verify the experimental findings, we have made ab initio *FP-LAPW* calculations of partial densities of states of Cd 4p, Ag 4p, and Ga 3d electrons that are presented in Figs. 4–6. In the theoretical data, the Fermi level is set to zero and the zero of energy in Figs. 4–6 has been taken at the top of the last occupied band as it was suggested in a number of theoretical ab initio band-structure calculations for semiconducting Cd- and Ga-containing sulfides and selenides [3, 6, 9]. It is worth mentioning that multi-peak structures of Cd 4p, Ag 4p, and Ga 3d states predicted by the *FP-LAPW* calculations for $AgCd_2GaS_4$ (Figs. 4–6) are not resolved by the XP data measured for this compound (Fig. 3) because of life-time and apparatus broadening the experimental spectra.

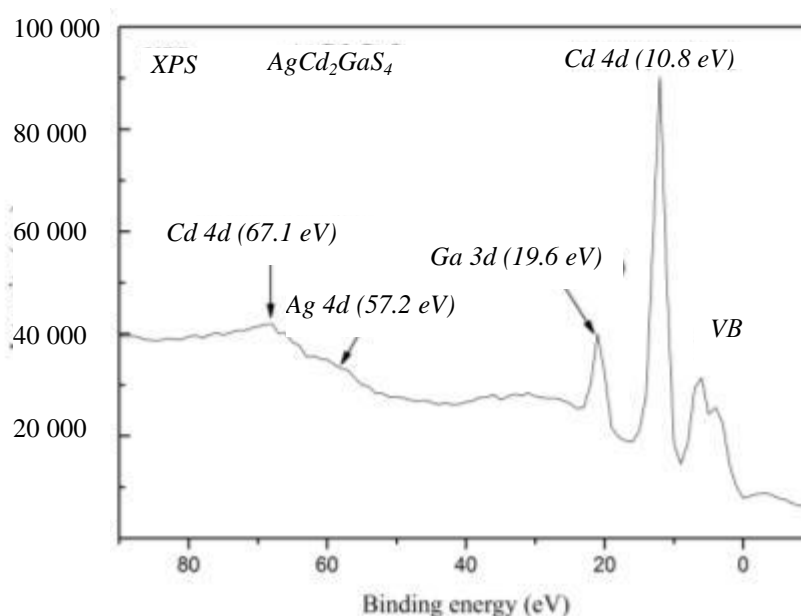


Fig. 3. XP valence-band spectrum including the energy region corresponding to location of upper core-levels of the atoms constituting a quaternary $AgCd_2GaS_4$ compound

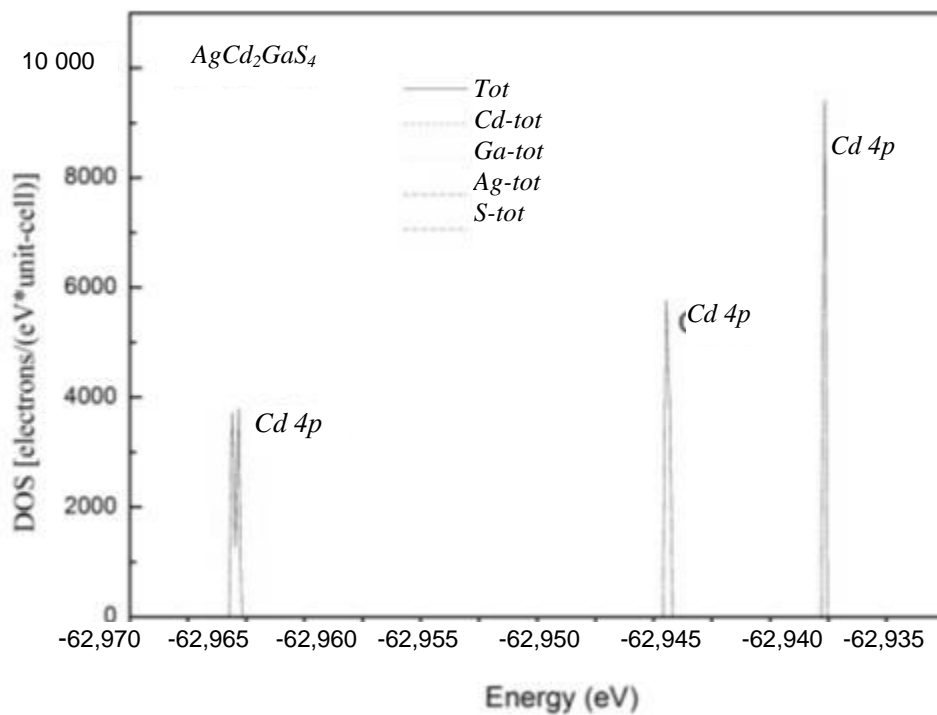


Fig. 4. Partial densities of the Cd 4p states of AgCd₂GaS₄

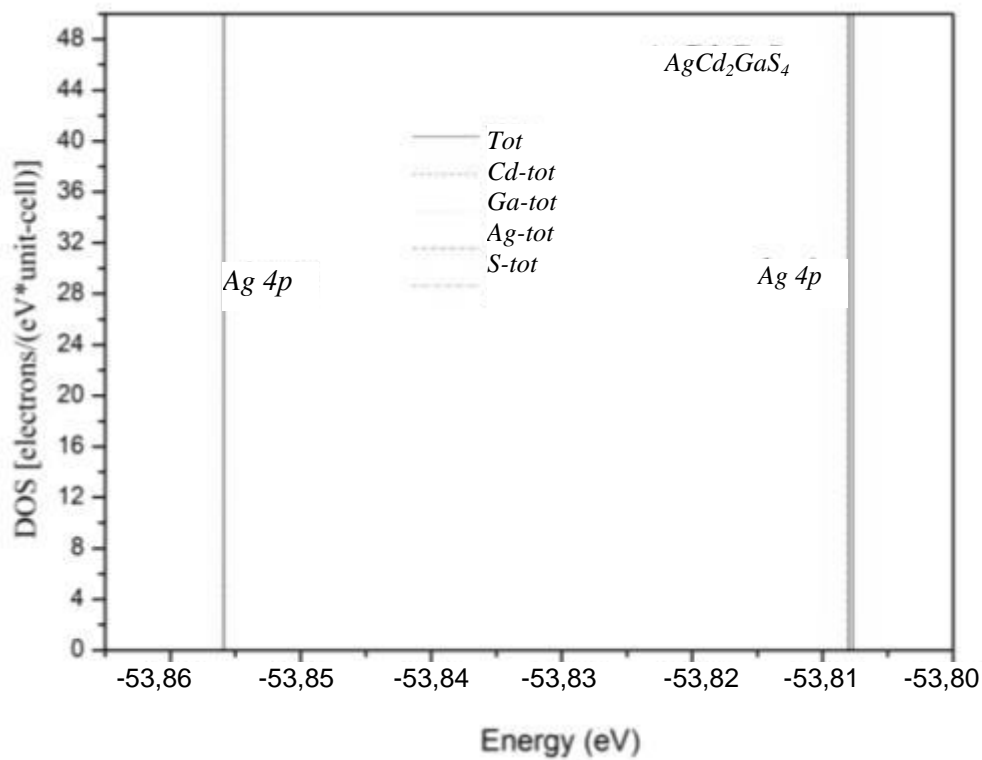


Fig. 5. Partial densities of the Ag 4p states of AgCd₂GaS₄

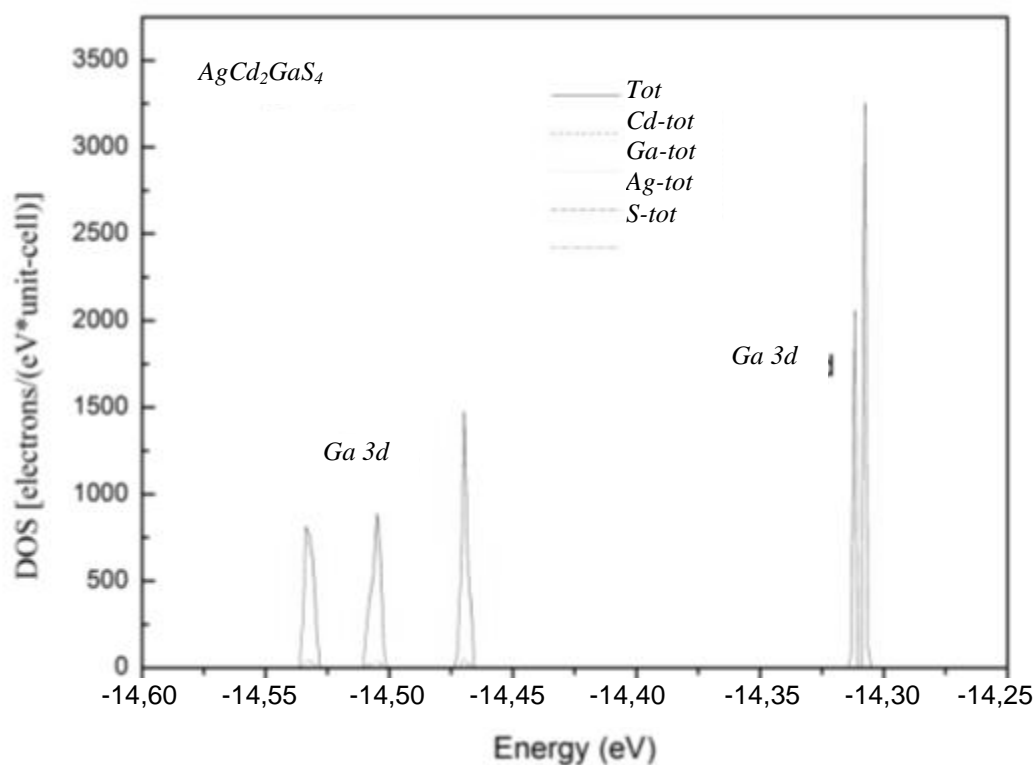


Fig. 6. Partial densities of the Ga 3d states of $\text{AgCd}_2\text{GaS}_4$

Fig. 7 shows comparison of the XP valence-band spectrum (shifted by 0,5 eV towards the Fermi energy) and the total DOS of the $\text{AgCd}_2\text{GaS}_4$ compound provided that a common energy scale is used. In addition, total DOS and the most important partial densities of states, mainly the Ag *d*, Cd *d*, Ga *p* and S *p* states, in the valence-band region of $\text{AgCd}_2\text{GaS}_4$ as established in Ref. [7] are also shown in Fig. 7, for comparison. The lowest band D ranging from $-7,6$ to $-8,2$ on the curve of total DOS of $\text{AgCd}_2\text{GaS}_4$ presented in Fig. 7 is formed almost exclusively due to contributions of the Cd *d* states. Some small contributions of the Ga *s* and S *p* states in the band D was also detected by the *FP-LAPW* calculations in Ref. [7]. As can be seen from Fig. 7, the curve of total DOS of $\text{AgCd}_2\text{GaS}_4$ contains the band C ranging from $-5,7$ to $-6,3$ eV, which is separated by a 1,3 eV gap from the band D. The band C is composed mainly due to contributions of the Ga *s* states. Minor contributions of the S *p* states in the band C were found to be characteristic of the electronic structure of the $\text{AgCd}_2\text{GaS}_4$ compound [7]. Furthermore, the sub-band B of the valence band of $\text{AgCd}_2\text{GaS}_4$ is composed mainly of the contributions of Ag *d* states. As one can see from Fig. 7, there exist a 0,8 eV gap between the band C and the sub-band B on the curve of total DOS of $\text{AgCd}_2\text{GaS}_4$. The upper sub-band A of the valence band of $\text{AgCd}_2\text{GaS}_4$ is composed mainly of contributions of the Ag *d* and S *p* states in almost equal proportions. Those states are strongly hybridized within the energy region corresponding to the position of the sub-band A of the curve of total DOS of $\text{AgCd}_2\text{GaS}_4$. As Fig. 7 demonstrates, energy positions of the sub-bands A and B of the theoretical curve of total DOS and the experimental XP valence-band spectrum of $\text{AgCd}_2\text{GaS}_4$ are close to each other. The binding energy position of the theoretical band D of $\text{AgCd}_2\text{GaS}_4$ is somewhat underestimated compared with the experimental XP measurements made for this compound.

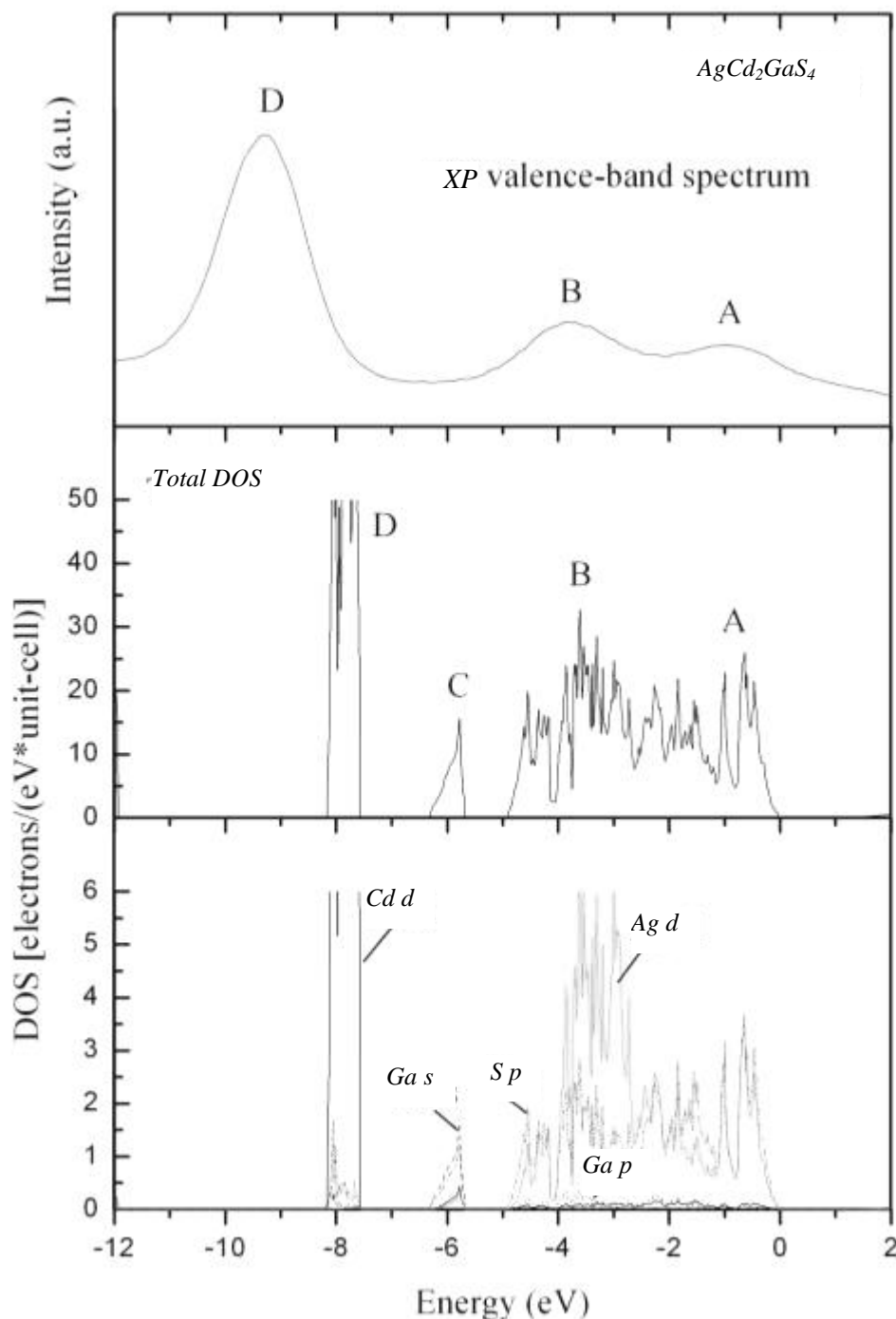


Fig. 7. Comparison on a common energy scale of curves of partial Ag d, Cd d, Ga s, and Sp densities of states (bottom panel), curve of total DOS (middle panel) and XP valence-band spectrum (upper panel) of $AgCd_2GaS_4$

It is worth mentioning that the band C presented on the curve of total DOS of $AgCd_2GaS_4$ is not visible on the XP valence-band spectrum of the compound as Fig. 7 reveals. This band is formed by contributions of the Ga s states as it has been mentioned above. This fact can be explained by minor photoionization cross-section value of Ga 4s electrons that is more than 25 times smaller compared to those of Ag 4d and Cd 4d electrons, the principal contributors to the sub-band B and the band D, respectively, on the curve of total DOS of $AgCd_2GaS_4$ (see Fig. 7) when exciting the XP spectra with a Mg K α source [16].

In addition, photoionization cross-section value for *Ga 4s* electrons is smaller more than one order of magnitude in comparison with that of *S 3p* electrons, other significant contributors to the sub-band A on the curve of total *DOS* of *AgCd₂GaS₄*.

Conclusions. The *XP* valence-band spectrum including the comparatively wide energy region corresponding to location of upper core-levels of the atoms constituting a quaternary *AgCd₂GaS₄* compound (up to ~ 80 eV with respect to the bottom of the valence band) has been measured. For such a purpose, *AgCd₂GaS₄* single crystal grown by the Bridgman method has been used. To verify the experimental data, electronic structure of *AgCd₂GaS₄* has been also theoretically studied employing the first-principles *FP-LAPW* method. With this aim, partial densities of states of upper core-levels for the *AgCd₂GaS₄* compound have been calculated. The theoretical and experimental data regarding energy positions of the *Cd 4p*, *Ag 4p* and *Ga 3d* core-level spectra of the *AgCd₂GaS₄* compound were found to be in agreement to each other.

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Божко Володимир, Третяк Аліна, Булатецька Леся, Бекеньов Валерій, Парасюк Олег, Хижун Олег. Електронна структура *AgCd₂GaS₄*. Рентгенівський фотоелектронний (РФ) спектр валентних електронів монокристалу *AgCd₂GaS₄*, вирощеного методом Бріджмена, досліджено в широкому енергетичному інтервалі, котрий уключав спектри внутрішніх електронів, розміщені на відстані аж до ~ 80 eV від нижнього краю

валентної зони. Електронна структура сполуки $AgCd_2GaS_4$ також розрахована «з перших принципів» методом приєднаних плоских хвиль – повного потенціалу (ППХ-ПП). Величини енергій зв'язку (E_{3b}) внутрішніх $Cd4p$ -, $Ag4p$ - і $Ga3d$ -електронів, що отримані теоретично для $AgCd_2GaS_4$ за допомогою ППХ-ПП методу, добре збігаються з експериментальними значеннями E_{3b} , котрі виміряні для досліджуваного кристала за допомогою РФ-спектроскопії. Виконано суміщення кривих парціальних щільностей станів складових атомів сполуки $AgCd_2GaS_4$ у єдиній енергетичній шкалі з РФ-спектром валентних електронів.

Ключові слова: $AgCd_2GaS_4$, електронна структура, розрахунок електронної структури з перших принципів, рентгенівська фотоелектронна спектроскопія, щільність станів.

Божко Володимир, Третяк Алина, Булатецкая Леся, Бекенцов Валерий, Парасюк Олег, Хижун Олег. **Электронная структура $AgCd_2GaS_4$.** Рентгеновский фотоэлектронный (РФ) спектр валентных электронов монокристалла $AgCd_2GaS_4$, выращенного методом Бриджмена, исследован в широком энергетическом интервале, который включал спектры внутренних электронов, расположенных на расстоянии до ~80 эВ от нижнего края валентной зоны. Электронная структура соединения $AgCd_2GaS_4$ также рассчитана «из первых принципов» методом присоединенных плоских волн – полного потенциала (ППВ-ПП). Величины энергий связи ($E_{св}$) внутренних $Cd4p$ -, $Ag4p$ - и $Ga3d$ -электронов, полученные теоретически для $AgCd_2GaS_4$ посредством ППВ-ПП метода, хорошо совпадают с экспериментальными значениями $E_{св}$, которые измерены для исследуемого кристалла с помощью РФ-спектроскопии. Выполнены совмещения кривых парциальных плотностей состояний составляющих атомов соединения $AgCd_2GaS_4$ в единой энергетической шкале с РФ-спектром валентных электронов.

Ключевые слова: $AgCd_2GaS_4$, электронная структура, расчет электронной структуры из первых принципов, рентгеновская фотоэлектронная спектроскопия, плотность состояний.

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Оптические, электрические та фотоэлектрические свойства кристаллов $Ag_2In_2Si(Ge)Se_6$

У роботі досліджено оптичні, електричні та фотоелектричні властивості маловивчених халькогенідних напівпровідників $Ag_2In_2Si(Ge)Se_6$. Показано, що досліджувані матеріали проявляють властивості неупорядкованих систем, обумовлених технологічними дефектами й структурними особливостями сполуки. Вивчено спектри оптичного поглинання напівпровідників $Ag_2In_2Si(Ge)Se_6$, розраховано коефіцієнт поглинання α , оцінено ширину забороненої зони. Інтерпретацію експериментальних результатів здійснено в межах моделі Мотта для неупорядкованих систем.

Ключові слова: напівпровідники, оптичне поглинання, дефекти, ширина забороненої зони.

Постановка наукової проблеми та її значення. Аналіз досліджень цієї проблеми. Кристали $AgInSe_2$, $AgInTe_2$ належать до великого класу з'єднань зі структурою халькопїриту $A^I B^{III} C_2^{VI}$, які мають технологічно привабливі властивості для створення різноманітних оптоелектронних приладів. Низка кристалів із цієї групи, такі як $CuInSe_2$, $CuGaSe_2$, $CuInS_2$, $AgGaS_2$, $AgGaSe_2$, є комерційно доступними й знайшли застосування у виробництві приладів нелінійної оптики та сонячних елементів. Останнім часом практичний інтерес мають й інші з'єднання цієї групи, у тому числі й $AgInSe_2$ та $AgInTe_2$, які є електронно-хімічними аналогами бінарних з'єднань $A^I B^{VI}$ [11].

Оптичні спектри вирощених у вигляді тонких плівок $AgInSe_2$ і $AgInTe_2$ вивчено в роботі [6]. Згідно з цими даними, кристали $AgInSe_2$ і $AgInTe_2$ є прямозонними й ширина забороненої зони становить 1,21 та 1,16еВ, відповідно, що робить їх перспективними матеріалами для перетворювачів сонячної енергії. Результати досліджень фотопровідності $AgInSe_2$ [8] засвідчили наявність хороших